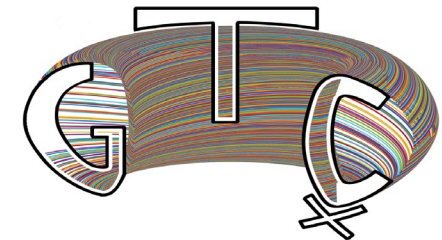
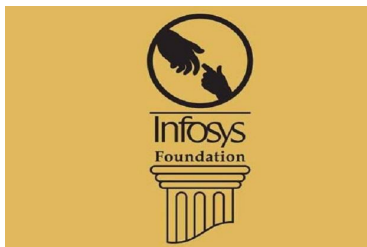


# Global gyrokinetic simulations of electrostatic microturbulent transport using kinetic electrons in LHD stellarator

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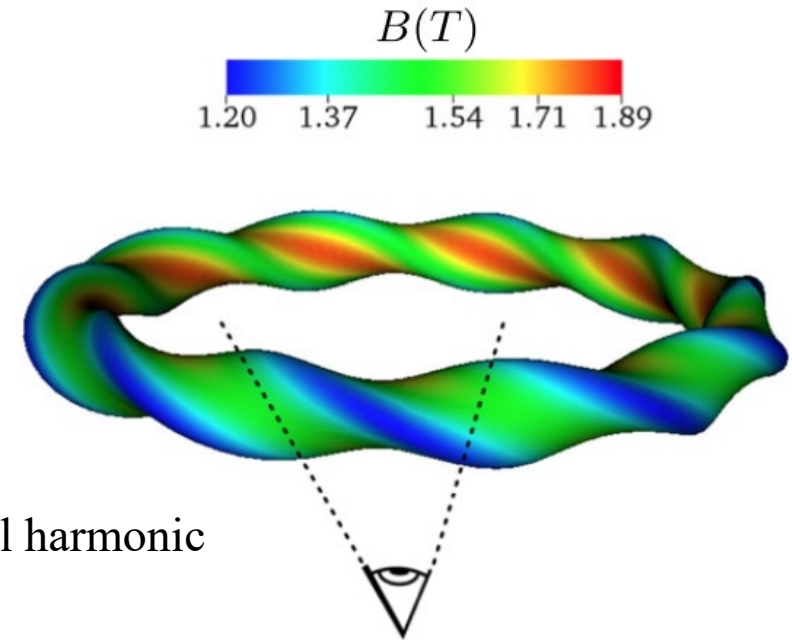
# LHD Equilibrium

- ❑ GTC uses non-axisymmetric equilibrium generated by ideal MHD code VMEC
- ❑ The equilibrium geometry and magnetic field are provided by VMEC as the Fourier series in both the poloidal and toroidal directions on a discrete radial mesh with equidistant in the toroidal flux.
- ❑ These 3D data in the left-handed straight fieldline coordinates are then transformed to the Boozer coordinates

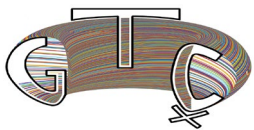
$$B(\psi, \theta, \zeta) = \sum_n [B_c(\psi, \theta, n)\cos(n\zeta) + B_s(\psi, \theta, n)\sin(n\zeta)]$$

n is the toroidal harmonic

- ❑ Field period in LHD is 10; therefore, all equilibrium quantities have periodicity  $2\pi/10$  in the toroidal direction
- ❑ Therefore, there are ten drift wave eigenmode families corresponding to the ten field periods for turbulent transport.



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# Simulation and Physics Model using GTC

## □ Dynamics in 5D phase space

$$\left( \partial_t + \dot{\vec{X}} \cdot \nabla + \dot{v}_{\parallel} \partial_{v_{\parallel}} \right) f_i(\vec{X}, \mu, v_{\parallel}, t) = C_i f_i,$$

$$\dot{\vec{X}} = v_{\parallel} \hat{b} + \vec{v}_E + \vec{v}_c + \vec{v}_g,$$

$$\dot{v}_{\parallel} = -\frac{\vec{B}^*}{m_i B} \cdot (\mu \nabla B + Z_i \nabla \phi),$$

Kinetic hybrid model for electrons

First order drift kinetic equation

$$\frac{dw_e}{dt} = \left( 1 - \frac{\delta f_e^{(0)}}{f_{0e}} - w_e \right) \left[ -\vec{v}_E \cdot \nabla \ln f_{0e}|_{v_{\perp}} - \frac{\partial}{\partial t} \left( \frac{e\delta\phi^{(0)}}{T_e} \right) - (\vec{v}_d + \delta\vec{v}_E) \cdot \nabla \left( \frac{e\phi}{T_e} \right) \right]$$

Lowest order (adiabatic)

$$\delta\phi = \delta\phi^{(0)} + \delta\phi^{(1)}$$

First-order (kinetic)

$$f_i = f_{0i} + \delta f_i$$

Lowest order (adiabatic)

$$f_e = f_{0e} + \delta f_e^{(0)} + \delta h_e \rightarrow \text{Higher order (kinetic)}$$

equilibrium

Gyrokinetic Poisson equation

$$\frac{e\tau}{T_e}(\phi - \tilde{\phi}) = \frac{\delta\bar{n}_i - \delta n_e}{n_0}$$

Lowest order

$$\frac{(\tau + 1)e\delta\phi^{(0)}}{T_e} - \frac{\tau e\delta\tilde{\phi}^{(0)}}{T_e} = \frac{\delta\bar{n}_i - \langle \delta\bar{n}_i \rangle}{n_0}$$

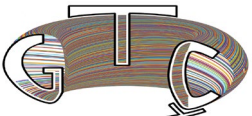
electrostatic potential to the first order is obtained from

$$e^{e\delta\phi/T_e} = e^{e\delta\phi^{(0)}/T_e} - \frac{\delta n_e - \langle \delta n_e \rangle}{n_0},$$

Higher order      Lowest order

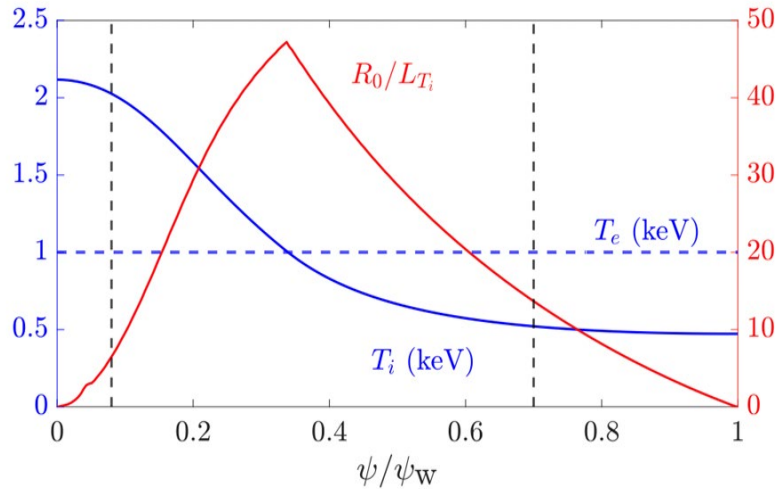
$$\delta n_e = \int \delta h_e d^3\vec{v}.$$

higher order electron density (kinetic)

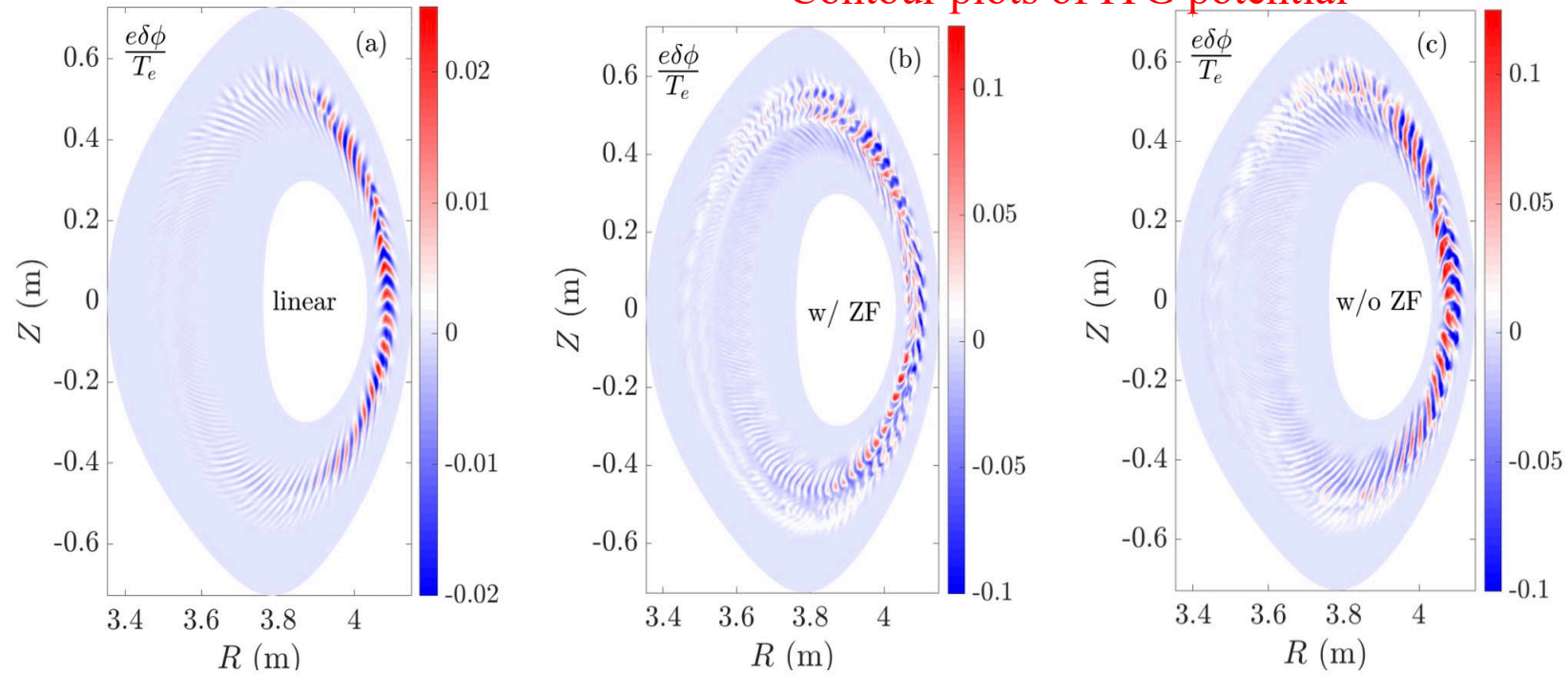


# Microinstabilities (ITG for $\eta = \infty$ ) in LHD

## Plasma profiles

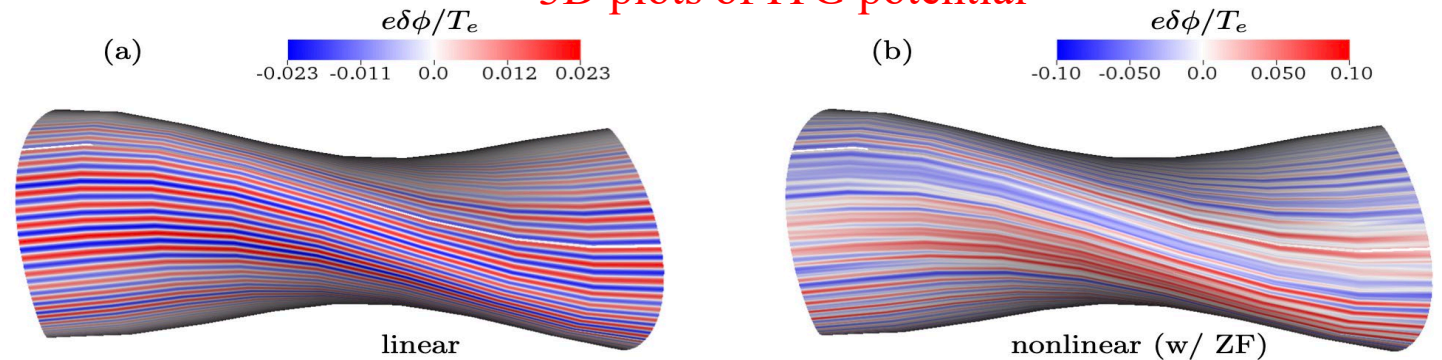


## Contour plots of ITG potential

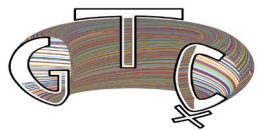


ITG mode structure is localized on the outer midplane, similar to tokamak

## 3D plots of ITG potential



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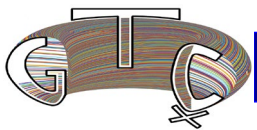
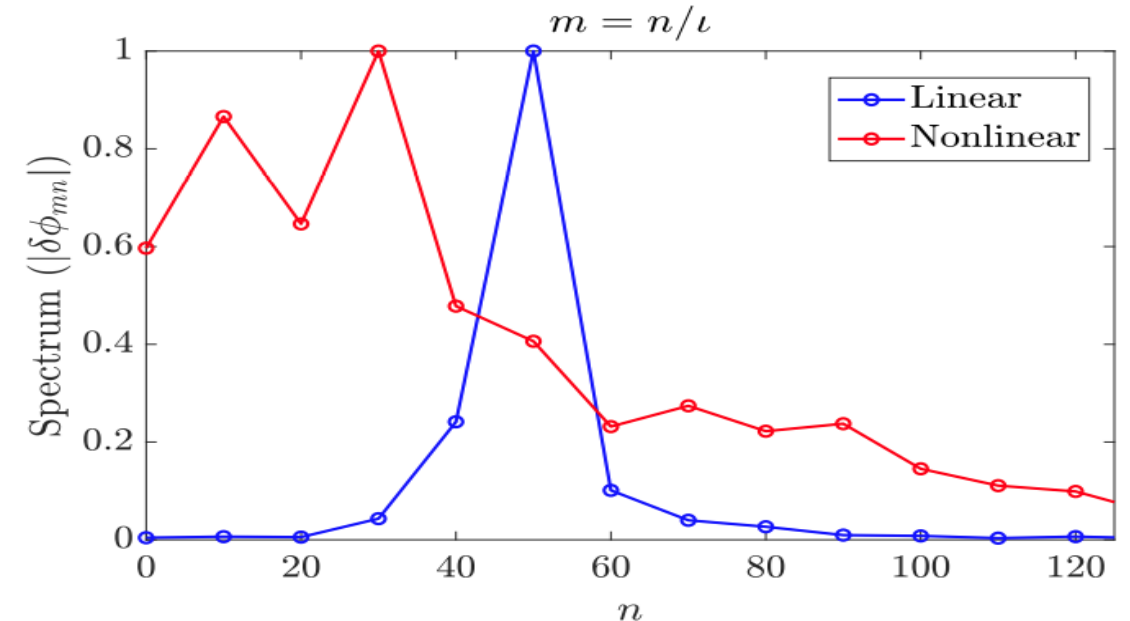
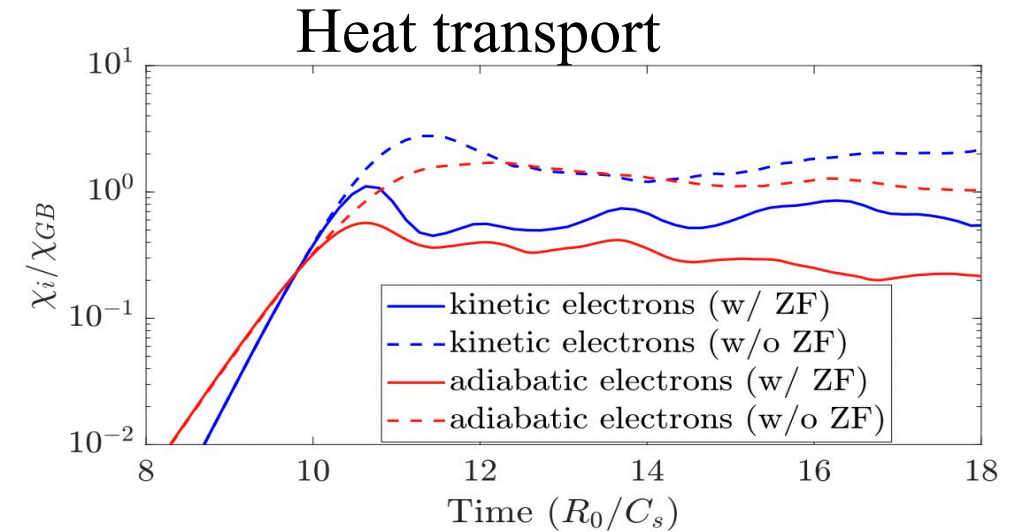


# Microturbulence (ITG for $\eta = \infty$ ) in LHD

- ❑ Kinetic electrons increase the ITG growth rate.
- ❑ Kinetic electrons reduce zonal flow and hence increase ITG turbulent transport.
- ❑ Zonal flow plays an important role in regulating the ITG plasma transport

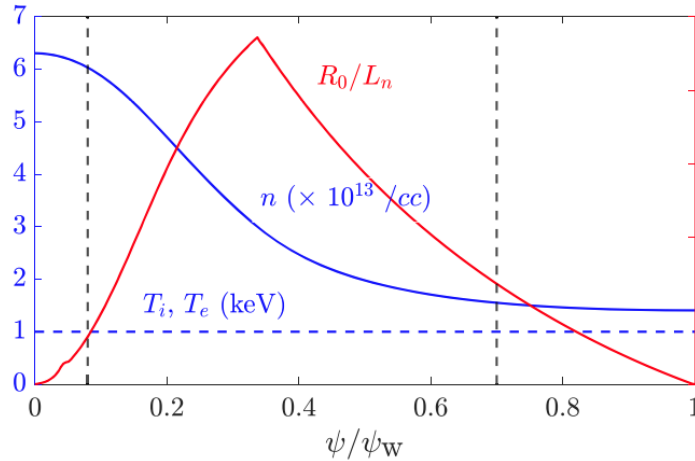
- ❑ The spectrum in the linear phase is narrow, width of  $\Delta n \sim 20$
- ❑ During the nonlinear phase, the toroidal spectrum after the inverse cascade becomes broader due to nonlinear mode coupling.

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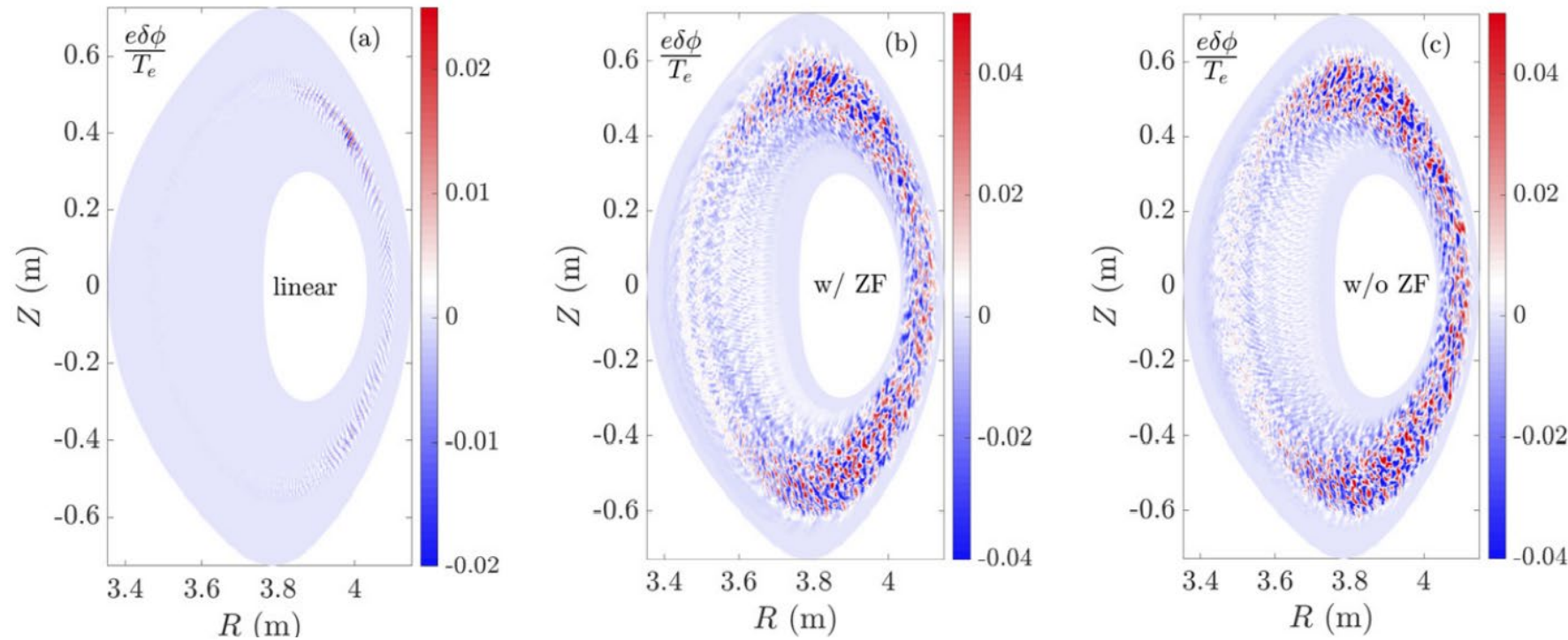


# Microinstabilities (TEM for $\eta = 0$ ) in LHD

## Plasma profiles

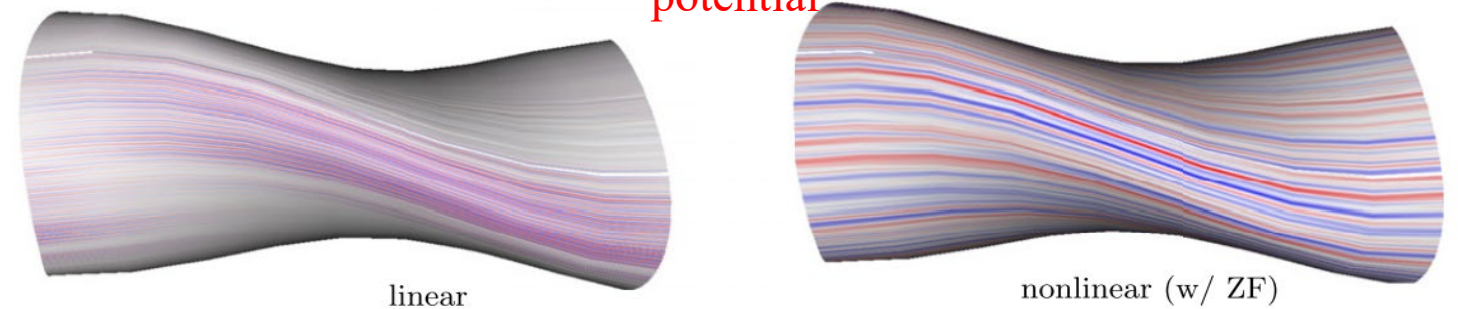


## Contour plots of TEM potential

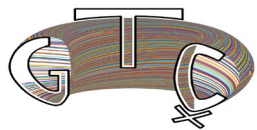


TEM mode structure is localized on the outer midplane, similar to tokamak

(a)  $e\delta\phi/T_e$  color scale: -0.023 to 0.023  
 (b)  $e\delta\phi/T_e$  color scale: -0.094 to 0.094  
 3D plots of TEM potential



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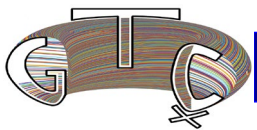
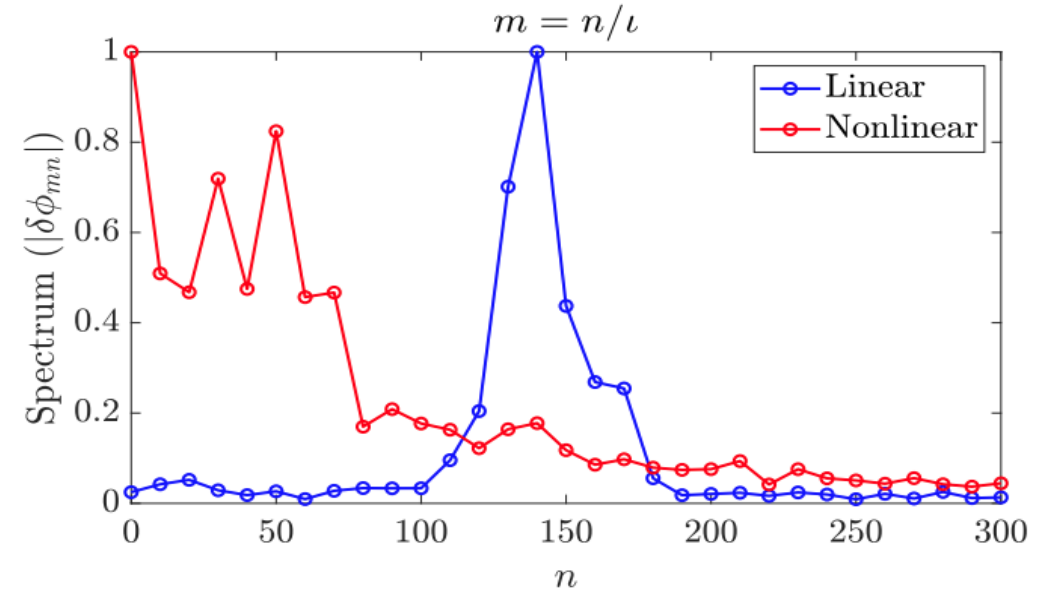
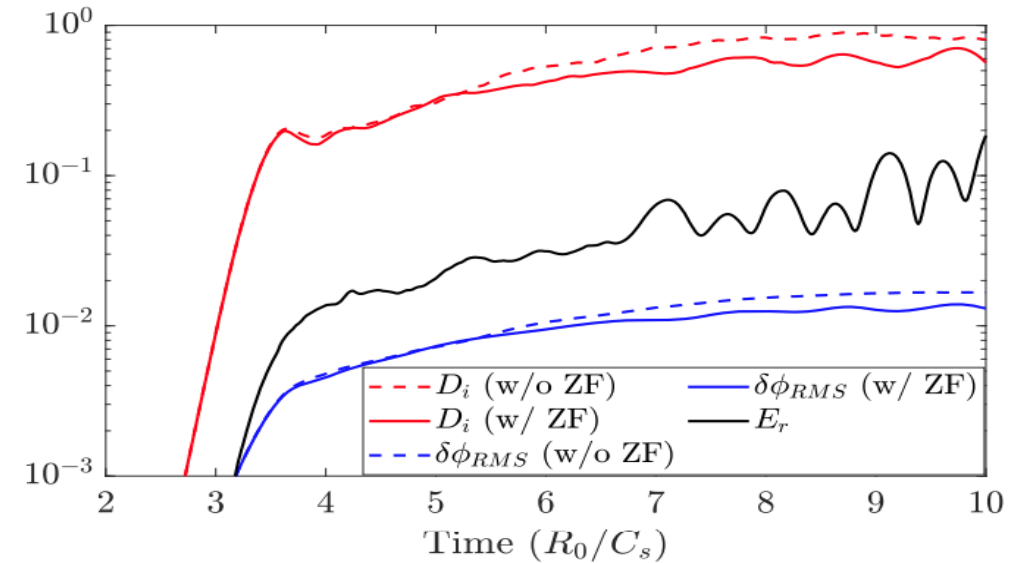
# Microturbulence (TEM for $\eta = 0$ ) in LHD

□ Turbulence spreading is observed during the nonlinear phase, but the turbulent eddies are barely affected by zonal flows. Unlike ITG turbulence, the zonal flow is not playing an important role in the case of TEM turbulence.

□ During the nonlinear saturation, the nonlinear poloidal and toroidal mode coupling leads to an inverse cascade from high to low mode numbers, the main reason for the saturation of the mechanism.

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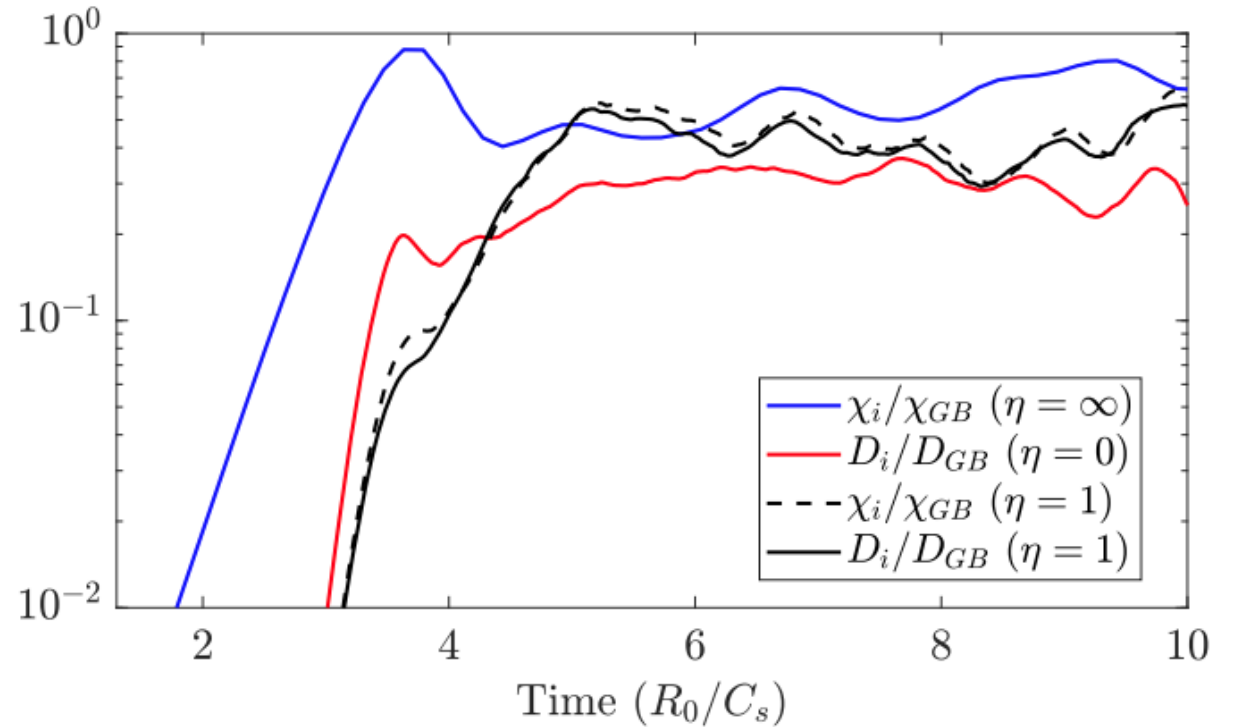
## Ion particle transport



# Microturbulence (for $\eta = 1$ ) in LHD

- ❑ The dominant mode is propagating in the ion diamagnetic direction: **ITG**
- ❑ The volume average ion diffusivity and conductivity are almost identical for  $\eta = 1$  case
- ❑ Transport with  $\eta = \infty$  case is the highest, while  $\eta = 0$  has the lowest transport level.

**Conclusion:** In the present scenario similar plasma profile gradients, the ITG turbulence acts as the primary drive for the heat conductivity transport, whereas the TEM turbulence is effective for the particle diffusivity



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